

COMMUNICATION

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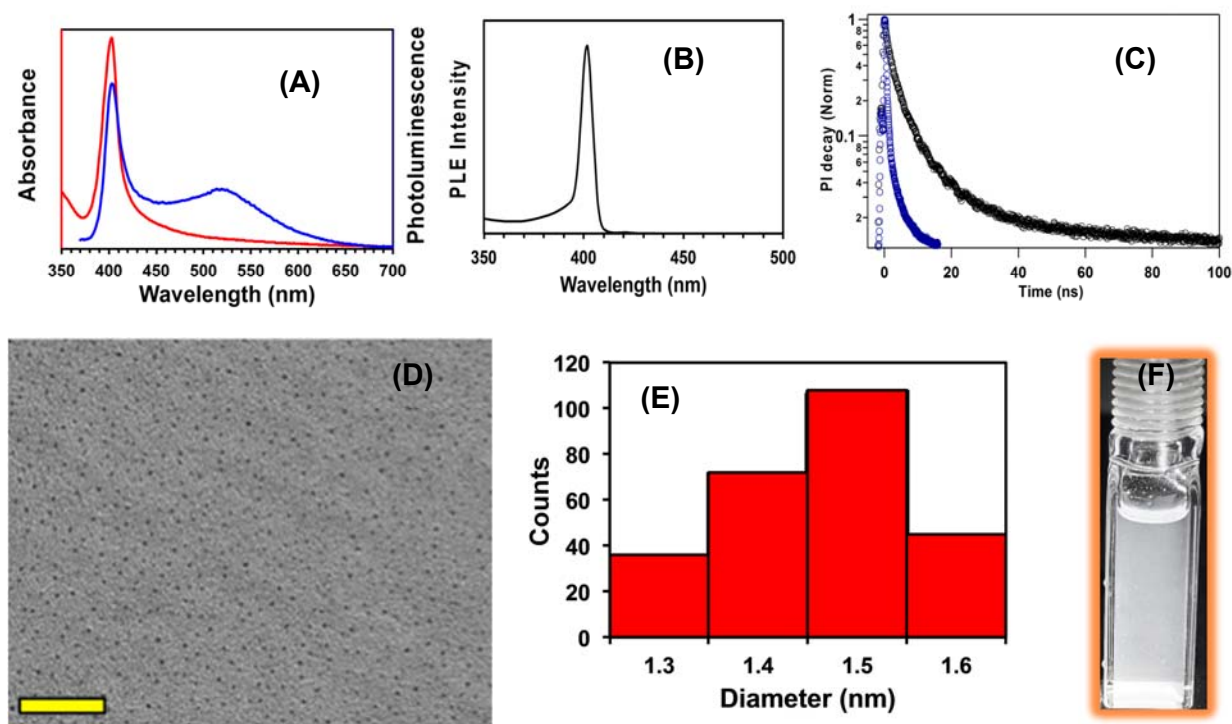


Fig. 1 (A) UV-visible absorption (red) and PL (blue) spectra of $\text{CH}_3\text{NH}_3\text{PbBr}_{3.0}$ PNCs at room temperature in toluene. The excitation wavelength was 350 nm. (B) PLE spectrum of $\text{CH}_3\text{NH}_3\text{PbBr}_{3.0}$ PNCs at room temperature in toluene. (C) PL decay traces (blue: band-edge at 409 nm; black: broadband at 515 nm) of the PNCs excited at 405 nm. (D) TEM image of the PNCs with 50 nm scale bar. (E) Histogram shows PNCs with an average 1.5 nm in diameter. (F) White-light emitting PNCs under illumination of UV light.

As illustrated in Fig. 1A, purified MAPbBr_3 PNCs in toluene display a sharp absorption peak at 402 nm. The PNCs exhibited a combination of band-edge (emission maxima at 403 nm) and broadband (emission maxima at 512 nm) PL emission that are markedly different than both organic-inorganic hybrid^{12–14} and inorganic^{15–18} perovskite nanocrystals, which display sharp band-edge emission features. The PLE spectrum (Fig. 1B) of the PNCs was identical in shape to that of the absorption spectrum. Therefore, our white-light PNCs have their lowest energy transition at 402 nm. The MAPbBr_3 PNCs display short emission lifetimes ($\tau_{\text{band-edge}}$: ~3 ns and $\tau_{\text{broadband}}$: 7 ns), see Fig. 1C. We hypothesize that the relatively longer broadband emission lifetime of our MAPbBr_3 PNCs in comparison to other perovskite quantum dots (QDs)¹⁹ is associated with a combination of multi-channel radiative recombination of excitons and delocalization of excitonic wave function, as discussed later. As shown in Fig. 1D, low-resolution transmission electron microscopy (TEM) analysis of the purified PNCs showed the presence of ultrasmall nanoclusters (see Fig. 1E) with a narrow size distribution. An average diameter of 1.5 nm (see Fig. 1E) was determined from the high-resolution TEM analysis (Fig. S1). The diameter of the PNCs corresponds to only ~2 unit cells. Therefore, it was extremely difficult to capture TEM images showing lattice spacing. Exposure (5 min) of the electron beam did not damage the PNCs. The PNCs CIE chromaticity coordinates were (0.304, 0.351) (Fig. S2), which fall within the white light region

of pure white light CIE coordinates of (0.333, 0.333) as perceived by the naked eye.

The sharp absorption feature indicates narrow size distribution of synthesized MAPbBr_3 PNCs that resemble ultrasmall CdSe nanoclusters (<2.0 nm in diameter).^{20–24} A narrow size distribution would result in a sharp band-edge PL peak rather than a combination of band-edge and broadband peaks. The broad PL feature and short PL lifetimes suggest the presence of surface related trap sites in the PNCs, not their size distribution. Our organic-inorganic hybrid PNCs are coated with primary amines, which are known to be weak π -donating, L-type ligands²⁵ and their interaction with Pb^{2+} in PNCs would lead to formation of antibonding orbitals below the conduction band of PNCs.²⁶ The newly formed orbitals serve as electron trapping mid-gap states if not fully passivated, resulting in broad emission in the PL spectra. Moreover, the broad-band emission dominates because the relative number of atoms on the surface of the PNCs increases as their size decreases.²⁷ In this context, ultrasmall CdSe nanoclusters are known to display a combination of band-edge and broadband PL emission due to large surface to volume ratio and the presence of mid-gap trap states.^{20,22,23,28,29} Therefore, the PL spectrum of our 1.5 nm diameter MAPbBr_3 PNCs nicely corroborate results from ultrasmall CdSe nanoclusters. We calculated PL quantum yield for the PNCs and found it to be 5%, which is at least 10-fold higher than the first generation, white-light emitting organic-lead halide bulk perovskites.⁹ However, quantum yield of our PNCs is much lower than

conventional 3.3-7.0 nm diameter MAPbBr₃ QDs (50-70%).^{12,30} It is known that the larger QDs display higher quantum yield than ultrasmall nanoclusters because of relatively fewer surface defects in the crystal structure of the QDs, which allows the majority of the exciton recombination to take place radiatively.

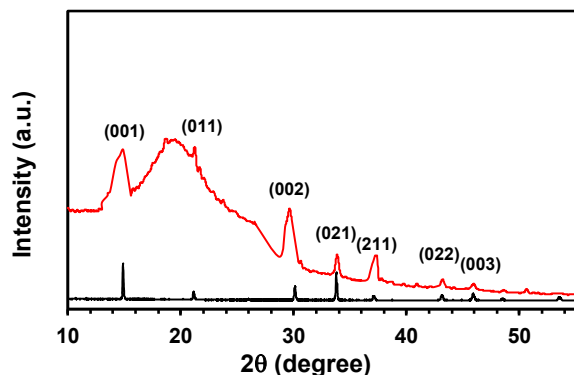


Fig. 2 X-ray diffraction (XRD) analysis of the bulk perovskite and ultrasmall PNCs. Observed XRD pattern of the bulk MAPbBr₃ (black) and purified ultrasmall MAPbBr₃ PNCs (red). The pattern of the PNCs corresponds to the cubic phase and the broadening of the diffraction peaks is characteristic of perovskite nanocrystals.^{12,15}

It is important to mention that our MAPbBr₃ PNCs displayed approximately 90 and 120 nm blue-shifts of band-gap in comparison to MAPbBr₃ QDs (3.3 nm diameter)¹² and bulk³¹ perovskite, respectively. Considering the relatively small exciton Bohr radius (~2.0 nm) of MAPbBr₃ (weak quantum confinement),³² the observed blue shifts of PNCs were large and significant. Recent effective mass approximation (EMA) calculations by our group³³ and others^{16,17,34} demonstrated that a simple spherical quantum well-based EMA calculation is not applicable to ultrasmall PNCs. A possible explanation is the potential delocalization of exciton wave functions in the strong confinement regime altering the confinement.

The energy dispersive spectroscopy (EDS) characterization (Fig. S3) showed bromide rich PNCs with Pb/Br atomic ratio of 3.8. The powder X-ray diffraction (XRD) analysis (red line) demonstrated broad peaks in comparison to the bulk perovskites (black line), two-dimensional (2D) materials,³⁵ and large anisotropically-shaped organolead halide perovskite nanostructures.³⁶ We should differentiate between our <2.0

nm diameter PNCs and 2D MAPbBr₃ material in terms of the Pb:Br ratio. It was reported that in 2D single-layered MAPbBr₃ perovskites, the Pb:Br ratio is 1:4 using a formula of [PbBr₄]²⁻.³⁵ The XRD pattern of our MAPbBr₃ PNCs is similar to a literature report of small MAPbBr₃ nanocrystals.¹² Our ultrasmall PNCs exhibited cubic structure (Fig. 2) and thus consist of octahedral unit cells of [PbBr₆]⁴⁻. Therefore, excess bromide ions should reside at the surface of the MAPbBr₃ PNCs. Our experimental data are in agreement with the literature suggesting an ~4.0 Pb/Br ratio for <2.0 nm diameter PNCs.¹² Therefore, because of higher bromide content in the structure of ultrasmall PNCs than traditional MAPbBr₃ perovskite, we refer to them as MAPbBr_{3.8} for the rest of this communication. Furthermore, the diameter of 1.34 nm calculated using Scherrer equation (see ESI file) is close proximity to the value characterized with TEM. Appearance of the broad peak at 1.52 ppm associated with the NH₂ resonance of amine in the ¹H NMR analysis confirmed the presence of a primary amine (hexylamine (HA) and/or diaminododecane (DADD)) at the surface of the PNCs (Fig. S4). Moreover, the α and β protons that were adjacent to -NH₂ appearing at 2.72 and 1.6 ppm, respectively, were also broad and resembled surface attached ligands.

Although absorption spectra and white light emission under photoexcitation of our ultrasmall MAPbBr₃ PNCs are similar to the literature reports of bulk 2D perovskite,^{9,10} atomically thin 2D perovskite,³⁵ and quasi-2D layered organic-PbBr₃ perovskite materials,³⁷ the physical, emission, and structural properties of our PNCs are markedly different than those 2D materials. (1) Our ultrasmall MAPbBr₃ PNCs are fully soluble in non-polar organic solvents such as toluene and chloroform (Fig. 1F), but 2D perovskite materials are not. Therefore, their optical characterizations of bulk and atomically thin 2D materials were conducted in solid-state^{9,10,35} as oppose to our colloidal-state analysis. (2) 2D organic-PbBr₃ perovskite display either sharp band-edge^{35,38} or broadband^{9,10} PL properties, but our ultrasmall MAPbBr_{3.8} PNCs display a combination of band-edge and broadband emission. (3) In addition, the diffraction peaks in our ultrasmall MAPbBr_{3.8} PNCs were dramatically broadened as compared to any-type of 2D perovskites.^{35,38} The XRD patterns of our PNCs are similar to the literature report of ~3.3 nm diameter MAPbBr₃ nanocrystals,¹² which display relatively sharper XRD peaks than our MAPbBr_{3.8} PNCs, suggesting an even smaller size than the literature report of 3.3 nm diameter MAPbBr₃

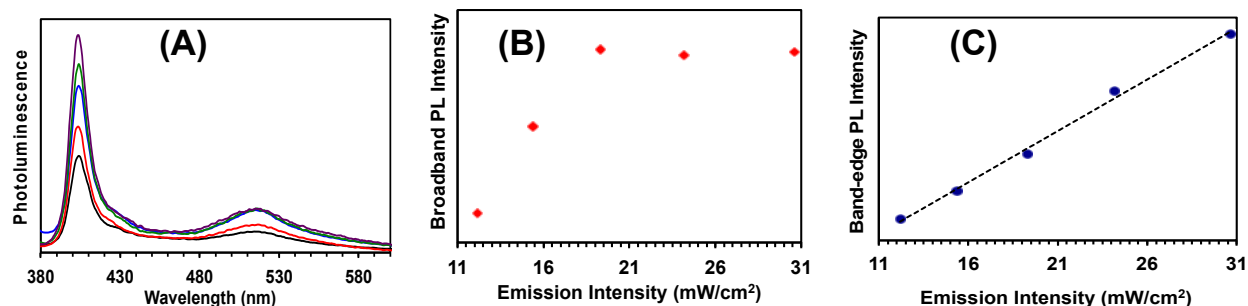


Fig. 3. (A) PL spectra of MAPbBr_{3.8} PNCs at different excitation power density (mW/cm²): black (11.2), red (15.4), blue (19.3), green (24.2), and purple (30.6). Dependence of broadband (B) and band-edge (C, R² = 0.995) PL intensity as a function of excitation power density of MAPbBr_{3.8} PNCs at room temperature.

nanocrystals. Finally, synthesis of bulk or atomically thin, 2D layered organic-PbBr₃ perovskite requires short chain diamines as surface passivating ligands.^{9,10,35} In this context, we have found that the presence of both medium chain length monoamine (HA) and long chain length diamine (DADD) in the reaction mixture is essential to obtain white-light emitting MAPbBr_{3.8} PNCs. In this context, synthesis of PNCs in the absence of DADD while keeping other conditions identical resulted in insoluble bulk MAPbBr₃ perovskite (Fig. S5). Furthermore, absence of HA produced white-light emitting MAPbBr₃ PNCs but QY was found to be <0.6% (see Fig. S6 for optical and XRD characterization). Therefore, the presence of HA and DADD onto the surface of MAPbBr₃ PNCs is a prerequisite to keep them in their ultrasmall size regime and maintain the surface related trap-states, resulting in white-light emission.

As shown in Fig. 1A, MAPbBr_{3.8} PNCs display a combination of band-edge and broadband PL properties. It is mentioned in the literature that the broad PL emission of 2D perovskite materials generally occurred from self-trapping of excitons or electrons/holes.^{10,39,40} This electronic process is different than our proposed mechanism, which involves radiative recombination of excitons at mid-gap trap-states giving our broad PL peak. To unravel the exact mechanism, we performed power dependent excitation PL studies (see Fig. 3A). Importantly, the broad PL emission peak became saturated as the power was increased (Fig. 3B). Thus, the experimental data do not support the self-trapping of excitons

because if this would occur the PL intensity would follow a linear relationship with excitation power.^{39,40} Considering that the trap-state concentration in MAPbBr_{3.8} PNCs is finite and slow to relax, trap-state PL would be expected to saturate at high excitation power, as shown for GaN.⁴¹ Therefore, our results strongly support trap-state-related emission. Importantly and as expected, the band-edge PL peak demonstrated a linear dependence with excitation power. Taken together, our experimental data suggest that the PL properties we observed (Fig. 1A) appear solely from the ultrasmall MAPbBr_{3.8} PNCs but not from a mixture of three-dimension (3D) and 2D structures. Furthermore, our optical data are in agreement with our XRD pattern that indicates ultrasmall PNCs of 3D crystal phase.

Based on the results described above, we propose that our white-light emission originates from contributions of sharp emission in the blue region and broad emission in the green region and thus is a combined effect from the large band-gap and presence of mid-gap trap states. In other words, by systematic manipulation of spectral features we should be able to control the PNCs CIE chromaticity coordinates as well as brighten the white-light emission. Therefore, our strategy to enhance these photophysical properties of MAPbBr_{3.8} PNCs involves increasing exciton binding energy and/or reducing trap states. To achieve this, we performed anion exchange reactions of purified, ultrasmall MAPbBr_{3.8} PNCs with methylammonium chloride (MACl). The EDS analysis (Fig. S7) confirmed the formation of mixed halide, MAPbCl_{1.3}Br_{2.5} PNCs,

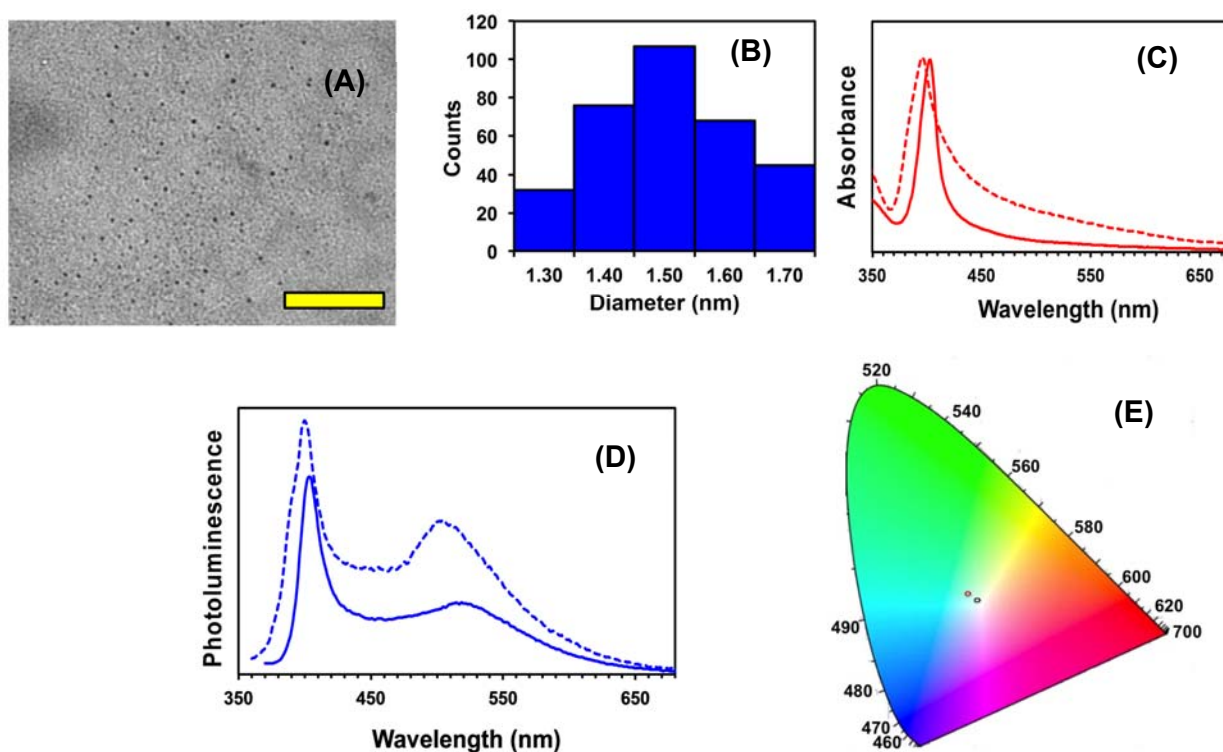


Fig. 4 (A) TEM image of the MAPbCl_{1.3}Br_{2.5} PNCs. The scale bar is 50 nm. (B) Histogram for PNCs shown in panel A. (C) UV-vis spectra of MAPbBr_{3.8} (solid line) and MAPbCl_{1.3}Br_{2.5} (dotted line) PNCs. (D) PL spectra of MAPbBr_{3.8} (solid line) and MAPbCl_{1.3}Br_{2.5} (dotted line) PNCs at 350 nm excitation. (E) CIE coordinates of MAPbBr_{3.8} (red) and MAPbCl_{1.3}Br_{2.5} (black) PNCs.

which displayed cubic crystal structure (Fig. S8). A slight shift of the (022) plane to a higher diffraction angle in MAPbCl_{1.3}Br_{2.5} PNCs in comparison to the original MAPbBr_{3.8} PNCs suggests exchange of bromide ions by chloride ions. Thus, we hypothesize that the majority of the surface bromide ions, along with some octahedral unit cells of [PbBr₆]⁴⁻ bromide ions, were exchanged with chloride ions (Fig. S9). TEM analysis confirmed no detectable change in the size of PNCs due to chloride exchange (Fig. 3A). However, due to their extremely small size, it would be difficult to accurately measure so slight a variation in the core diameter of PNCs.

Absorption measurement of our MAPbCl_{1.3}Br_{2.5} PNCs showed an ~6 nm blue shift in the band-gap (Fig. 4C), which is in agreement with the EDS analysis described above and the literature that chloride substitution of both bulk organic-PbBr₃ perovskites⁴² and MAPbBr₃ nanocrystals^{43,44} increases the band-gap because of the stabilization of their valence band with more electronegative chloride as compared to bromide.⁴⁵ Small blue-shift is also in agreement with the suggestion that fewer octahedral unit cells of [PbBr₆]⁴⁻ bromide ions were exchanged with chloride ions. Thus, the cubic crystal structure of the PNCs is dominated by the bromide ions. In the PL characterization (Fig. 4D), we observed changes in the position and shape of the broad-band emission in mixed halide PNCs that could suggest an influence of chloride in the excited states, which are involved in the emission properties.¹⁰ Moreover, increase in the band-gap increases the exciton binding energy⁴⁶ that results in an enhancement of the possibility of more efficient recombination of the excitons.²⁷ We observed a QY up to 12% of MAPbCl_{1.3}Br_{2.5} and radiative lifetimes ($\tau_{\text{band-edge}}$: ~5 ns and $\tau_{\text{broadband}}$: 19 ns) see Fig. S10. The broadband emission lifetime is more than two fold higher than MAPbBr_{3.8} PNCs. Moreover, we observed nearly homogeneous emission (determined from the PL lifetime) of our MAPbCl_{1.3}Br_{2.5} PNCs (data not shown). Importantly, the PLQY of our MAPbCl_{1.3}Br_{2.5} nanoclusters is higher than the white-light emitting bulk organic-lead halide perovskites (PLQY = 9%).¹⁰ The CIE chromaticity coordinates (Fig. 4E) of MAPbCl_{1.3}Br_{2.5} are (0.332, 0.341) which are very close to the CIE coordinates of pure white-light. Perhaps as recently demonstrated for CsPbBr₃ nanocrystals,^{16,47} control halide exchange and preparation of white-light emitting inorganic

perovskite nanoclusters are feasible if one were able to synthesize them in their ultrasmall size regime. We should mention that prolonged chloride exchange of MAPbBr_{3.8} PNCs resulted in insoluble white materials, thus demonstrating the importance of controlling the structural modification of nanoclusters to obtain desired photophysical properties.

Our structural and spectroscopic characterizations suggest that the bromide at the surface of the MAPbBr_{3.8} PNCs were exchanged with chloride. Because of the higher electronegativity of chloride than bromide, surface exchange reduces the surface related radiative mid-gap trap states⁴⁸ and increases the QY and radiative lifetimes as observed previously for CdSe²⁸ and CdTe nanoclusters,⁴⁹ respectively. In order to further investigate the trap state-related emission, we studied the excitation wavelength dependent PL properties. Fig. 5A demonstrates the broadband emission intensity as a function of excitation wavelength of MAPbBr_{3.8} and MAPbCl_{1.3}Br_{2.5} PNCs. The nonlinearity in the relationship along with change in the broadband emission peak in MAPbBr_{3.8} PNCs suggest the presence of variable defect sites,⁵⁰ which were not observed for MAPbCl_{1.3}Br_{2.5} PNCs as a consequence of the reduction of nonradiative surface mid-gap trap states. We also observed the change in the broadband peak shape between these two PNCs (Fig. 5B), suggesting the presence of different trap sites. Another important observation we made is the increase in $\tau_{\text{broadband}}$ of MAPbCl_{1.3}Br_{2.5} over that of MAPbBr_{3.8} PNCs. However, one should expect the opposite effect in lifetime after chloride substitution because of the increase in confinement energy that would result in higher exciton binding energy and efficient recombination of excitons. It is reported that the kinetic energy of excitons increases as the confinement energy increases.^{46,51} Under such circumstance, the excitonic wave functions could easily fill the organic-inorganic hybrid core of MAPbCl_{1.3}Br_{2.5} more completely than MAPbBr_{3.8} PNCs, and thus be prone to delocalize into the ligand monolayer. The delocalization of exciton wave functions increases the PL lifetime and PLQY, as recently reported for ultrasmall CdSe nanoclusters.⁵² Nevertheless, further investigation is required to fully delineate the origin of the white-light emission of these ultrasmall organic-inorganic PNCs. The unique photophysical properties of these PNCs have tremendous potential for SSL applications, as we

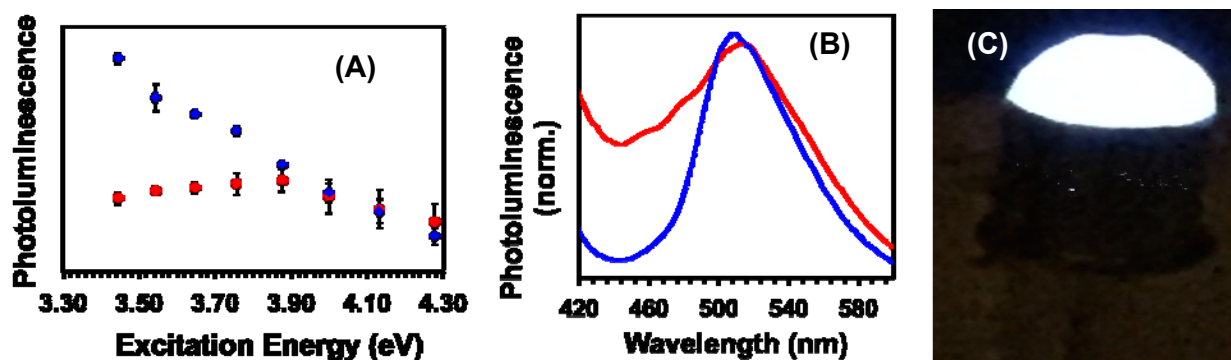


Fig. 5 (A) Excitation wavelength dependent broadband PL intensity of MAPbBr_{3.8} (red) and MAPbCl_{1.3}Br_{2.5} (blue) PNCs. (B) Comparison of broadband PL peak of MAPbBr_{3.8} (red) and MAPbCl_{1.3}Br_{2.5} (blue) PNCs at 350 nm excitation. (C) A 5 mm commercial UV LED (360 nm) coated with MAPbCl_{1.3}Br_{2.5} PNCs in polyurethane.

demonstrated with a proof-of-concept white-light source fabrication using a thin coating of MAPbCl_{1.3}Br_{2.5} PNCs in polyurethane (Fig. 5C).

Recently, it was reported that perovskite QDs display excellent PLQY,^{12,17-19,53} high carrier mobilities, and larger diffusion lengths,⁵⁴ which are crucial electronic properties for photovoltaic applications. These electronic properties arise because of excellent passivation of the QDs surface and reduction of the nonradiative trap states. However, the very narrow band-edge emission characteristic of QDs makes it difficult to produce white light emission and fabricate LEDs for SSL applications. In this present study, we utilized surface-related trap states (defect sites) to generate a complimentary broadband emission in PNCs. The combined band-edge and broadband emissions of our PNCs cover most of the visible region of the solar spectrum, which in turn results in white-light emission. Taken together, manipulation of the appropriate amount of radiative trap states in PNCs will enhance their optoelectronic properties and facilitate their applications in white light emitting LED fabrication.

In summary, we have developed the first hybrid perovskite nanocluster synthetic method that is capable of producing nearly pure white light, thus enhancing the possibility to use these PNCs as single source materials for SSL applications. Our PNCs would provide three important beneficial features in LED fabrication in comparison to traditionally used phosphors (organic and/or inorganic) or mixtures of different light-emitting QDs: (1) Mixing of multiple single light-emitting phosphors is commonly performed to cover the entire visible region of the solar spectrum. This mixing process reduces the device efficiency because the different phosphors decompose at different rates. (2) Inorganic phosphors require metal ion doping in order to achieve white-light emission.^{55,56} (3) White-light emitting multi-layer films can be prepared by mixing various color-emitting QDs but the devices suffer significant deficiency due to the self-absorption process that arises from different sized nanoclusters. Our organic-inorganic hybrid PNCs obviate all these drawbacks. Moreover, these PNCs have the potential to provide fundamentally new optical and electronic properties at the molecule-nanocrystal interface that could be utilized to maximize charge transport efficiency in optoelectronic devices.

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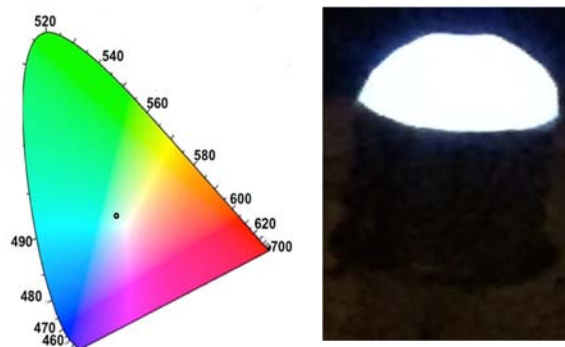
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TOC



Pure white-light emitting ultrasmall organic-inorganic hybrid perovskite nanoclusters were synthesized through a simple colloidal approach. Their emission properties were controlled by manipulating the surface chemistry.